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Calculating the Envelope Correlation Coefficient Directly From Spherical Modes Spectrum

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Abstract—This paper proposes an improved technique to calculate the Envelope Correlation Coefficient (ECC) by directly using the spherical mode spectrum of the antenna under test. The technique avoids errors due to numerical pattern integration and simplifies post-processing for near-field measurements. The technique is successfully tested on two different antenna types, in the total number of four MIMO configurations.

Index Terms—MIMO, envelope correlation coefficient, spherical near field, antenna measurement.

I. INTRODUCTION

The constantly growing need for higher wireless data throughput has resulted in common use of MIMO technology across the globe. To support this trend there is a great and increasing number of MIMO antenna designs, many of which are incorporated into a small volume [1-4]. Such antennas need to be characterized with respect to their performance within MIMO systems and considering often complex interactions between multiple antenna ports.

The Envelope Correlation Coefficient (ECC) has become a standard parameter for MIMO antennas and is defined as:

\[
ρ_{ECC} = \frac{\iint_{4\pi} E_1(\theta, \phi) \cdot E_2(\theta, \phi) d\Omega}{\iint_{4\pi} E_1(\theta, \phi) \cdot E_1(\theta, \phi) d\Omega \iint_{4\pi} E_2(\theta, \phi) \cdot E_2(\theta, \phi) d\Omega}
\]  

(1)

where \(E_1\) and \(E_2\) are far-field radiation patterns, generated from ports 1 and 2 respectively. The metric can be seen as a correlation of two radiation patterns, which, under the assumption that the incoming signals are isotropically distributed, yields also the correlation of MIMO signals at receiver inputs [5]. The lower the correlation, the higher data throughput can be supported by the MIMO antenna. To calculate (1), one needs to perform a full sphere antenna measurement for each port.

Hallbjörner in [6] has proposed a simplified approach to approximate the ECC when only S-parameters and efficiencies \(\eta_1\) and \(\eta_2\) are known:

\[
ρ_{ECC} = -\frac{S_{11}^* S_{12} + S_{21}^* S_{22}}{\sqrt{(1 - |S_{11}|^2)(1 - |S_{21}|^2)(1 - |S_{22}|^2 - |S_{12}|^2 \eta_1 \eta_2)}
\]

(2)

This formula is valid if the losses between both antennas are either independent or if the loss correlation is small enough to be neglected. Although measurement of S-parameters is straightforward, measurement of efficiencies poses certain challenges. A technique to measure efficiencies of MIMO antennas with Wheeler’s Cap has been recently proposed by Moharram and Kishk in [7] but provides only limited accuracy. In most cases, measurements in an anechoic chamber are required for accurate estimation of efficiencies, thus limiting the benefits of (2). Please also note that neglecting the efficiency term in (2) will lead to a large error in the calculated ECC [8]. Thus to determine the ECC an anechoic chamber or reverberation chamber is required.

Spherical Near-Field (SNF) measurements are recently becoming more popular, especially for characterization of small antennas. Rather than directly sampling far-field radiation patterns, the technique calculates spherical modes of the antenna under test from near-field measurements [9]. The far-field radiation pattern is then calculated from the modal spectrum. The number of spherical modes is proportional to \((r^{1/2}/\lambda)^2\), where \(r\) is a radius of a sphere that fully encloses the Antenna Under Test (AUT).

In this paper we demonstrate, that the ECC can be calculated directly from the spherical mode spectrum of an antenna, without the need for far-field patterns. This approach requires less post-processing and offers greater precision, as it avoids the error when the radiation pattern is numerically integrated over a sphere. Furthermore, it provides a more intuitive understanding of the relation between antenna size and possible number of uncorrelated data streams.

II. PRINCIPLES

The spherical far-field antenna radiation pattern can be described as a superposition of spherical modes:

\[
E(\theta, \phi) = k_0 \sum_{l=1}^{\infty} \sum_{m=-l}^{l} \sum_{s=1}^{2} q_{sm} k_{sm}(\theta, \phi)
\]

(3)

where \(K_{sm}\) denotes the spherical harmonics, \(q_{sm}\) its complex coefficient, \(k_0\) is the wavenumber and \(Z_0\) is the impedance of free space. The order \(m\) \((-l \leq m \leq l)\) of the spherical harmonics describes the azimuthal variation \(\phi\) of the field. The variation in \(\theta\) depends on the degree \(l\) and the order \(m\). The index \(s \in \{1,2\}\) is connected with the components of transversal electric (TE) and transversal magnetic (TM) waves. The upper index \(c\) specifies the radial function of the spherical wave [9].
By definition the modes are orthogonal to each other, thus their correlation can be expressed using Dirac’s delta:

\[
\int_{4\pi} K_{s_{\mu}l_{\sigma}}^{(3)}(\theta, \phi) \cdot K_{s_{\mu}l_{\sigma}}^{(4)}(\theta, \phi) d\Omega = \delta_{s_{\sigma}l_{\nu}} \delta_{s_{\mu}l_{\nu}} (-1)^m
\]

with

\[
\sum_{\mu=1}^{\infty} \sum_{\nu=1}^{\infty} \sum_{\sigma=1}^{2} q_{s_{\mu}l_{\sigma}} K_{s_{\mu}l_{\sigma}}^{(3)}(\theta, \phi) = \sum_{\nu=1}^{\infty} \sum_{\mu=1}^{\infty} \sum_{\sigma=1}^{2} (-1)^m(q_{\sigma,-\nu,\mu})^2 K_{s_{\mu}l_{\sigma}}^{(4)}(\theta, \phi)
\]

and by combining (4) and (3) into (1) we obtain:

\[
\rho_{ECC} = \frac{\sum_{\nu=1}^{\infty} \sum_{\mu=1}^{\infty} \sum_{\sigma=1}^{2} q_{s_{\mu}l_{\sigma}}^2(q_{\nu,-\sigma,\mu})^2}{(\sum_{\nu=1}^{\infty} \sum_{\mu=1}^{\infty} \sum_{\sigma=1}^{2} q_{s_{\mu}l_{\sigma}}^2)(\sum_{\nu=1}^{\infty} \sum_{\mu=1}^{\infty} \sum_{\sigma=1}^{2} q_{\nu,-\sigma,\mu}^2)}
\]

This allows the calculation of the ECC directly from the mode spectrum, i.e. the coefficients \(q_{s\mu l\nu}\). Please note, that the field of each mode \(K_{s_{\mu}l_{\sigma}}\) is irrelevant and only the coefficients contribute to the ECC value.

### III. Measurements and Discussion

Eq. (5) was tested with measured data from three two-element Printed Inverted-F Antennas (PIFA) and one dual-mode patch antenna. The measurements were performed in spherical near-field antenna measurement facility of the Institute of High Frequency Technology, RWTH Aachen University.

#### A. Printed Inverted F Antennas

A simple PIFA antenna was designed in [10] to operate in the bandwidth of 2.45 - 2.6 GHz. To validate (5) three different configurations were prepared as a testbed, each consisting of two PIFAs. The configurations varied both orientation and distance between antennas in order to provide representative span of possible ECC values. This involved an opposite configuration, with antennas placed back-to-back at a distance of 48 mm (Fig. 1a); an orthogonal configuration with 30 mm spacing between antenna feeds (Fig. 1b); and parallel configuration with 20 mm spacing between antenna feeds (Fig. 1c).

![Fig. 1. Three various PIFA configuration used as a testbed: a) opposite with 48 mm distance between feeds; b) orthogonal with 30 mm distance between feeds; c) parallel with 20 mm distance between feeds.](image)

#### B. Multimode patch

Multimode patch antennas have recently grown in popularity, as they offer a MIMO performance within a cost-efficient, low profile structure [3-4]. The antenna used as a testbed operates at 5.05 GHz and consists of two stacked patches, operating in TM\(_{10}\) and TM\(_{40}\) mode. Fig. 4 shows the ECC values calculated with (1) and (5) and again the curves are very close to each other.

![Fig. 2. ECC values calculated from (5) and (1) for the three investigated PIFA configurations.](image)

![Fig. 3. Absolute ECC differences calculated by (5) and (1) for the three investigated PIFA configurations.](image)
The differences for the examples are caused by a numerical integration error when using the standard trapezoidal method. We used the trapz function from Matlab for the spherical numerical integration via the trapezoidal method. In order to investigate this error in more detail, a perfect uncorrelated (ECC = 0) antenna configuration was simulated. The antenna consists of a horizontal Hertzian dipole (port 1) and two vertical Hertzian dipoles port 2a/b separated by $\lambda/2$ (see Fig. 5). Port 2b has a 180° phase shift compared to port 2a creating a null in the xy-plane. The radiated modes are listed in Table 1 and the radiation pattern should be perfectly uncorrelated. However, the calculated ECC according to (1) is affected by numerical errors and approaches zero with increasing angular resolution (see Fig. 6). Please note, that this error will change, depending on the spherical modes excited by each antenna and differs with antenna topology, as seen in Sections III.A and III.B.

The ECC based on the spherical mode spectrum does not have this dependency and is thus not affected by the numerical integration error.

### Table 1. $|q_{sim}|$ of Simulated Antenna Configuration

<table>
<thead>
<tr>
<th>Port 1</th>
<th>Port 2</th>
</tr>
</thead>
<tbody>
<tr>
<td>Mode $(s, m, l)$</td>
<td>$</td>
</tr>
<tr>
<td>$(2, -1,1)$</td>
<td>$\sqrt{2}$</td>
</tr>
<tr>
<td>$(2, +1,1)$</td>
<td>$\sqrt{2}$</td>
</tr>
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<td></td>
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### IV. Conclusions

The paper introduces a method to calculate the Envelope Correlation Coefficient directly from the spherical mode spectrum, without the need to extract the radiation pattern. The method was validated with four exemplary antennas. Furthermore, it was demonstrated that the proposed technique avoids the typical numerical integration error.

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### References


